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Paradamite

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Abstract. Zn_2AsO_4OH , P1, Z = 2, $D_c = 4.595$ Mg m⁻³; cell constants at 293 K: a = 5.638 (3), b = 5.827 (3), c = 6.692 (2) Å, a = 103.25 (4), $\beta = 104.37$ (3), $\gamma = 87.72$ (4)°; λ (Mo Ka) = 0.71069 Å. Paradamite is isostructural with tarbuttite, Zn_2PO_4OH .

Introduction. Paradamite is a dimorph of adamite and was first characterized by Switzer (1956). Finney (1966) derived the unit cell and proposed that paradamite is isostructural with tarbuttite, Zn_2PO_4OH , and Hawthorne (1976) has discussed the relationship between the adamite and tarbuttite structure types.

Single-crystal precession photographs of paradamite from Mapimi, Mexico, exhibit triclinic symmetry; the unit cell chosen is analogous to that used by Cocco, Fanfani & Zanazzi (1966) for tarbuttite. The unit-cell parameters were derived by least-squares refinement of 15 high-angle reflexions aligned automatically on a four-circle diffractometer. A full set of intensity data (1384 reflexions) was collected over one hemisphere out to a 2θ of 65° from an irregular equidimensional crystal of radius 0.09 mm, following the experimental procedure of Hawthorne & Ferguson (1975). Standard data-reduction procedures with a spherical absorption correction ($\mu = 20.2 \text{ mm}^{-1}$) resulted in 1202 unique reflexions of which 1132 were considered as observed $(I_{net} > 3\sigma)$. Structure refinement was carried out on an IBM 370/158 computer using the program RFINE (Finger, 1969). Scattering factors for neutral atoms were taken from Cromer & Mann (1968) with anomalous-dispersion corrections from Cromer & Liberman (1970). The atomic parameters given by Cocco *et al.* (1966) for tarbuttite, Zn_2PO_4OH , were assumed as the starting parameters for paradamite. Several cycles of full-matrix least-squares refinement resulted in convergence (for isotropic temperature factors) at a conventional R of 6.2% (observed reflexions). Conversion to anisotropic temperature factors of the form $exp(-\sum_{i=1}^{3}\sum_{j=1}^{3}h_ih_j\beta_{ij})$ resulted in convergence (maximum shift/error <0.01) at an R of 4.8% and an R_w (unit weights) of 5.8%. Final atomic positions and equivalent isotropic temperature factors are given in Table 1. Interatomic distances and angles

 Table 1. Atomic positions and equivalent isotropic temperature factors (Å²)

	x	У	Z	B_{eq}
Zn(1)	0.3928 (2)	0.2498 (2)	0.5008 (2)	0.86 (2)
Zn(2)	0.0214(2)	0.7355 (2)	0.1867 (1)	0.72 (2)
As	0.8312(2)	0.2476(1)	0.2745(1)	0.49 (2)
O(1)	0.7627 (11)	0.9689 (10)	0.1391 (9)	0.79 (9)
O(2)	0.9421 (12)	0.3885 (11)	0.1233 (9)	0.94 (9)
O(3)	0.5779 (11)	0.3837(11)	0.3342(9)	0.85 (9)
O(4)	0.9549 (11)	0.7478 (11)	0.4968 (9)	0.71 (8)
OH	0-3519 (11)	0.8899 (12)	0.2997 (9)	0.89 (9)

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were calculated with the program *ERRORS* (L. W. Finger, personal communication) and are given in Table 2.*

* A list of structure factors and details of the anisotropic thermal model have been deposited with the British Library Lending Division as Supplementary Publication No. SUP 34097 (7 pp.). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

Table 2. Selected interatomic distances (Å) and
angles (°) for paradamite

Zn(1)–O(3)	1.994 (6)	Zn(2)O(1)	1.972 (6)
Zn(1)O(3)	2.160 (6)	Zn(2)-O(2)	2.013 (6)
Zn(1)-O(4)	1.965 (6)	Zn(2) - O(2)	2.093 (6)
Zn(1)–OH	2.203 (6)	Zn(2)-O(4)	2.183 (6)
Zn(1)–OH	1.992 (6)	Zn(2)–OH	1.994 (6)
$\langle Zn(1)-O \rangle$	2.063	$\langle Zn(2)-O \rangle$	2.051
AsO(1)	1.676 (6)		
As-O(2)	1.680 (6)		
As-O(3)	1.696 (6)		
As-O(4)	1.693 (6)		
⟨As−O⟩	1.686		
Zn(1) trigonal	bipyramid		
O(3)–O(4)	3.627 (9)	O(3) - Zn(1) - O(4)	132.7 (3)
O(3)-OH	3.154 (9)	O(3)-Zn(1)-OH	104.6 (3)
O(4)OH	3.470 (9)	O(4)-Zn(1)-OH	122.6 (3)
$\langle O - O \rangle^m$	3.417	$\langle O-Zn(1)-O\rangle^m$	120.0
$\dot{O}(3) - O(3)^{s}$	2.650(12)	O(3) - Zn(1) - O(3)	79.2 (3)
O(3) - OH	3.129(9)	O(3) - Zn(1) - OH	96.3 (2)
O(3) - O(4)	2.891(9)	O(3) - Zn(1) - O(4)	88.9 (2)
O(3) - OH	3.200(9)	O(3) - Zn(1) - OH	100.8 (3)
O(4) - OH	2.978(9)	O(4) - Zn(1) - OH	91.0 (2)
OH-OH'	2.828(12)	OH-Zn(1)-OH	84.6 (2)
$\langle 0 - 0 \rangle^a$	2.946	$\langle \mathbf{O} - \mathbf{Z} \mathbf{n}(1) - \mathbf{O} \rangle^a$	90.1
	O(3) - Zn(1) - 0	OH 173.6 (2)	
7n(2) trigonal	binyramid		
$\Sigma II(2) II Igoliai$	2 477 (0)	O(1) = 7r(2) = O(2)	121 5 (2)
O(1) - O(2)	3.477(9)	O(1) - Zn(2) - O(2)	121.5(3)
O(1) - OH	3.278 (9)	O(1) - ZII(2) - OH O(2) - Zn(2) - OH	117.1(3)
O(2) = OH	3.307 (9)	$(2) = 2 \ln(2) = 0 \ln(2)$	127.1 (5)
$\langle 0 = 0 \rangle^m$	3.447	$(0-2\pi(2)-0)^{-1}$	00.0(2)
O(1) - O(2)	$3 \cdot 113(9)$	O(1) - Zn(2) - O(2) O(1) - Zn(2) - O(4)	99.9 (3)
O(1) = O(4)	2.920(0)	O(1) - Zn(2) - O(4)	76.2 (3)
$O(2) = O(2)^{2}$	2.355(12)	O(2) = Zn(2) = O(2) O(2) = 7n(2) = O(4)	85.7 (2)
O(2) - O(4)	2.020 (0)	O(2) = Zn(2) = O(4)	95.6 (3)
O(2) = OH	3.089 (9)	O(2) = Zn(2) = OH O(4) = 7n(2) = OH	95.3 (2)
(4) = 011	2 0 2 5	$(-7)^{-1} (2)^{-1} (3)^{-1} $	00.4
{ U - U }-	2.923	$(4) = \frac{161.9}{2}$	30.4
	0(2) - 2n(2) - 0	(4) 101.9 (2)	
As tetrahedror	1		100.0 (2)
O(1) - O(2)	2.715 (8)	O(1) - As - O(2)	108.0 (3
O(1) - O(3)	2·/08 (8)	O(1) - AS - O(3)	110.2 (2)
O(1) - O(4)	2.104 (8)	O(1) - AS - O(4)	110.2 (2)
O(2) - O(3)	2.107 (7)	$O(2) - A_{3} - O(3)$	100.5 (2)
O(2) - O(4)	2.751 (8)	$O(2) = A_3 = O(4)$ $O(3) = A_3 = O(4)$	108.5 (3)
(3) - 0(4)	2.131 (0)	(3)	100.5 (5
(1) = (1)	1.1.74	$\langle U - AS - U \rangle$	103.2

(m) Meridional parameters. (a) Axial parameters. (s) Shared edges.

Discussion. Cocco et al. (1966) indicate that the H atom is strongly bonded to the O atom designated OH. Difference Fourier maps in the vicinity of the OH anion showed a slight rather smeared peak between OH and O(1) that was assumed to be the H atom. However, the H position was not incorporated into the refinement as H positions derived from X-ray crystal structure refinement are usually significantly in error due to bonding effects. Examination of the empirical bond-strength table for paradamite (Table 3) shows a bond-strength excess around OH and a bond-strength deficiency around O(1). This together with the OH-O(1) distance of 2.877 (9) Å suggests significant hydrogen bonding between OH and O(1) (Fig. 1). Comparison of the bond-strength tables for adamite (Hawthorne, 1976) and paradamite suggests that the hydrogen bonding in paradamite is significantly stronger than in adamite.

Unfortunately there was insufficient material available to record an infrared spectrum of paradamite. However, comparison of the principal OH stretching frequencies in adamite (358.0 mm^{-1}) and tarbuttite (344.0 mm^{-1}) indicates that hydrogen bonding is

Table 3. Empirical bond-strength table for paradamite

Calculated from the curves of Brown & Shannon (1973).

	Zn(1)	Zn(2)	As	Н	Σ	Σ*
O(1)		0.469	1.265	(0.213)	1.734	(1.947)
O(2)		0∙423 0∙348	1.253		2.024	
O(3)	{0.444 0.298		1.205		1.947	
O(4)	0.478	0.282	1.214		1.974	
ОН	(0·270 (0·446	0.444	(0.787)	1.160	(1.947)	
Σ	1.936	1.966	4.937	(1.000)		

* Hydrogen-bond strengths estimated (values in parentheses) by minimization of deviation from ideality of sums around the anions.



Fig. 1. Polyhedral representation of the paradamite structure. The hydrogen bond is indicated by the broken line.

stronger in tarbuttite than in adamite, and lends support to the contention that it is also stronger in paradamite than in adamite.

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Refinement of 3C Pyrrhotite, Fe₇S₈

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Abstract. A metastable form of Fe_7S_8 , trigonal, $P3_121$, a = 6.8652 (6), c = 17.046 (2) Å, Z = 3, $D_x = 4.59$ Mg m⁻³. The crystal structure was refined to R =0.043 for 1026 reflections. The apparent space group $P6_222$ is explained by the twin-related domains with trigonal symmetry. The supercell contains ordered vacancies in alternating Fe layers normal to the c direction. The Fe atom is octahedrally coordinated with a mean Fe–S distance of 2.445 Å.

Introduction. Fleet (1971) determined the structure of $3C \text{ Fe}_7\text{S}_8$ assuming the twin-related domains with space group $P3_1$. The structural study of synthetic $3C \text{ Fe}_7\text{S}_8$ was made by Nishiguchi (1977), also based on $P3_1$. Their structures are essentially the same as that of $3C \text{ Fe}_7\text{S}_8$ reported by Okazaki & Hirakawa (1956) which has the symmetry $P3_121$. However, Fleet and Nishiguchi did not examine the possibility of $P3_121$. This study has been carried out to obtain the true symmetry and to determine precisely the actual displacements of atoms from the ideal structure.

The 3C Fe₇S₈ was prepared by the dry method and quenched from 973 K into ice water. A spherical crystal, 140 μ m in diameter, was used for data collection on an automatic four-circle diffractometer, using the ω -2 θ scan technique ($2\theta_{max} = 65^{\circ}$) with Mo

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 K_{α} radiation ($\lambda = 0.71069$ Å), monochromatized by pyrolytic graphite. The observed Laue symmetry (6/mmm) and the systematic absences (00l with $l \neq 3n$) indicate an apparent space group $P6_222$ or $P6_422$. Because the ordered 3C type of Okazaki & Hirakawa's (1956) model must have a trigonal symmetry of $P3_121$ ($P3_221$) or $P3_1$ (or $P3_2$), the apparent $P6_222$ (and/or $P6_422$) symmetry of the specimen has been explained by many small domains related by 180° rotation about [001] and/or by reflection across (100) in the specimen.

Observed structure factors, $F_o(hkl)$, were calculated from the observed intensities, $I_o(hkl)$, by the following relation:

$$|F_o(hkl)|^2 = (Lp)^{-1} AE_s I_o(hkl)I_c(hkl)/[I_c(hkl)]$$

+ I_c(khl)].

 $F_o(hkl)$ were corrected for the usual Lorentz and polarization factors (Lp), absorption $[A, \mu(Mo K\alpha) =$ 12.8 mm⁻¹, transmission factors between 0.27 and 0.33] and isotropic secondary extinction $[E_s, c =$ 0.59 (1) × 10⁻⁵; Zachariasen, 1963]. $I_c(hkl)$ and $I_c(khl)$ were calculated at each cycle of the leastsquares refinements, and the computing program *ORFLS* (Busing, Martin & Levy, 1962) was modified for this purpose. The scattering factors for neutral © 1979 International Union of Crystallography